

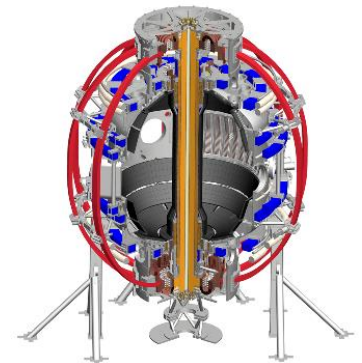
# Large Scale Data Analysis Case Study: NSTX-U Langmuir Probes

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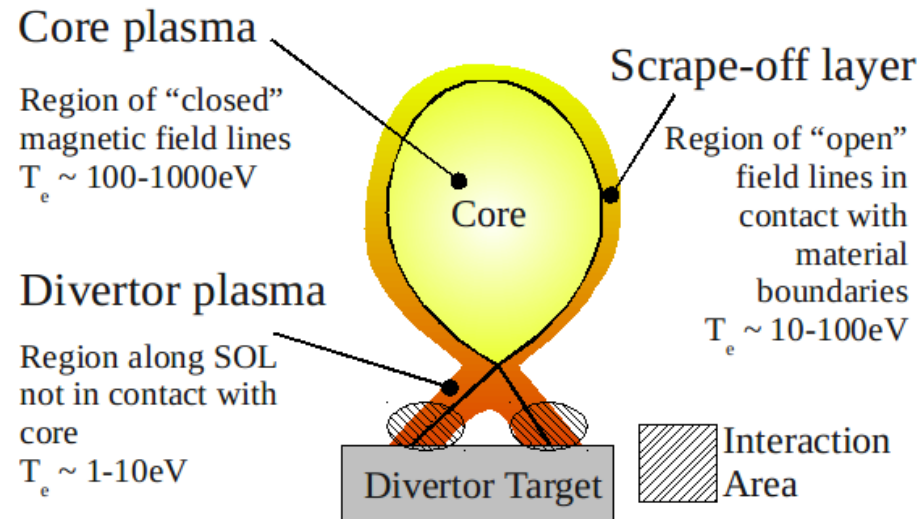
# Outline of Material

- Will you have too much data and why that means you have to think
- Strategies for attacking your data
- Langmuir Probes as a case study of where things fall apart
- Extracting Electron Distribution Functions: practical aspects and recommended practices

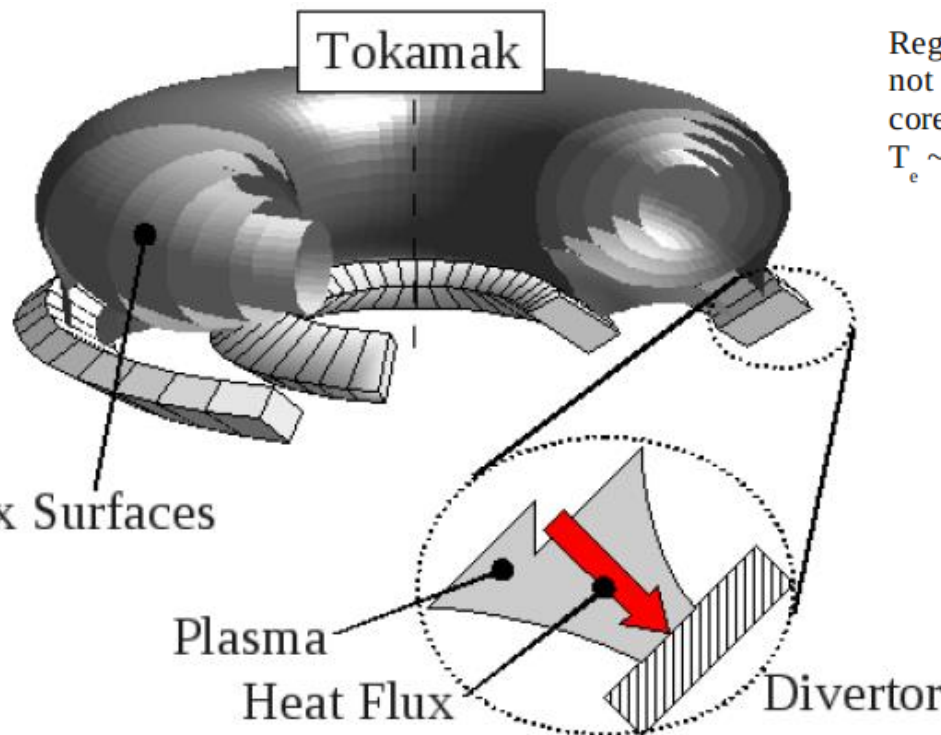
# Diverted tokamaks are the most developed concept for MFE

- Magnetic fields confine plasma in toroidal geometry
- Divertors developed to separate eroded material from core plasma

## A Diverted Tokamak

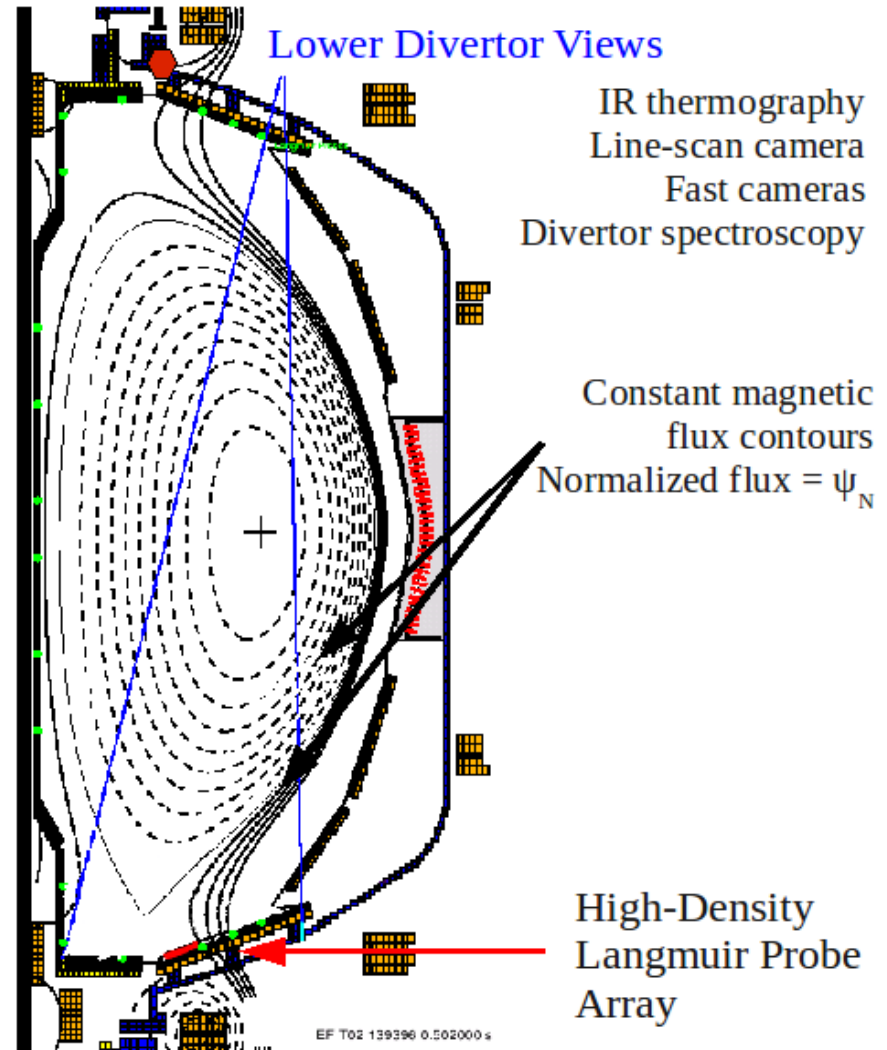
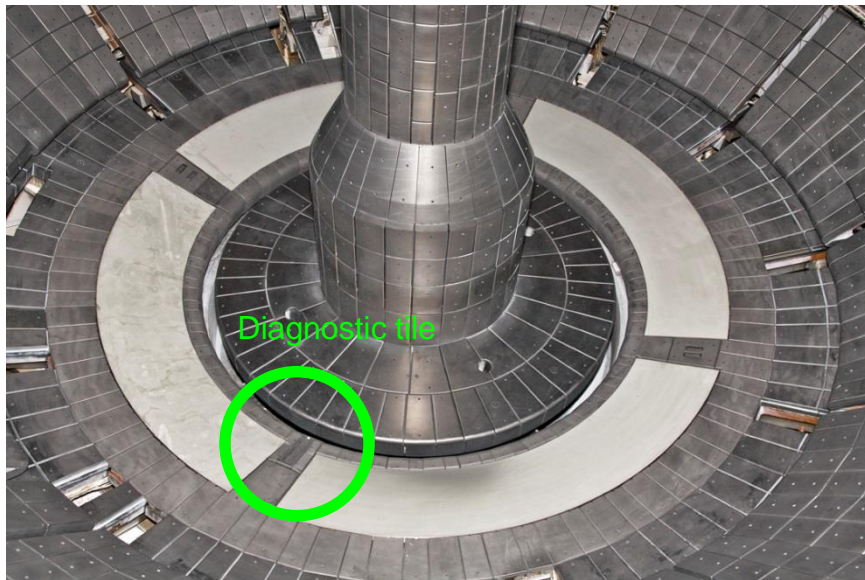


- A lot of processes happen in dense region at target plate – we need to diagnose it!



# Array of plasma diagnostics utilized during experiments

- Local plasma conditions and PFC currents measured with Langmuir probe system
- 2D fast-cameras provide nearly full toroidal coverage of divertor



# How much data is this?

- Each I-V characteristic can be a single “data set”
  - 250 kHz sampling frequency on digitizer
  - 500 Hz sweep rate, 2 V-sweeps per cycle
  - 1000 sweeps per 1s discharge per probe
  - 2000 shots per year
- **$2000 * 1000 * 4 = 8$  Million data sets in 1 year**
- What are the bigger picture questions and strategies needed to pare this down?

# What if I just analyze all the data by hand?

- Just not practical, but even if it were: would you want to?
- Scientific reporting is to provide sufficient information to replicate the work
  - Hand analysis means \*you\* are making decisions during analysis
  - If \*you\* are integral to analysis, no one else can repeat it!
- Systematic analysis allows study of the method itself
- Avoids “cherry-picking” if its systematic

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# Strategies for data analysis

1. Know what you want to get from the data!
  - E.g. I want the most accurate Ne and Te in divertor
  - E.g. I want the post-ELM evolution in the divertor
2. Know what you don't want in your data sets!
  - E.g. I don't want to analyze turbulent plasmas
  - E.g. I \*do\* want to analyze turbulent plasmas
3. Be Selective using 1 and 2!
  - Large data sets means you can “afford” to throw some things away
  - Don't slow yourself down analyzing every single case
4. Build in consistency checks (recommended practice)
  - Ensure your systematic analysis is being correctly implemented



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# Langmuir probes as case study

- The (almost) original plasma diagnostic (I. Langmuir 1881-1957)
- “Simple to implement; hard to interpret”
- Integrated measure of charged particles in plasma
- LP theory is deep and complex, lots of regimes of operation (see Demidov 2002)
- Perturbative diagnostic (at least locally)



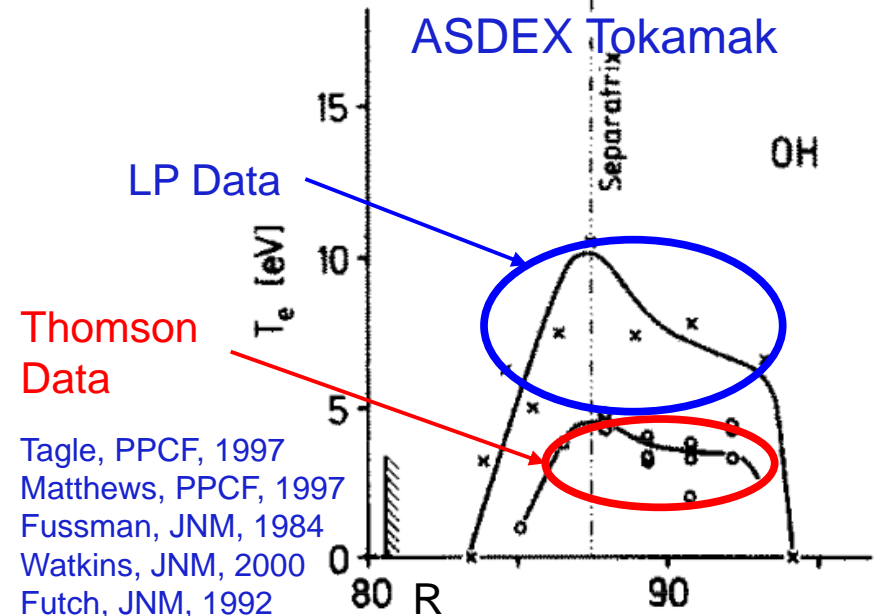
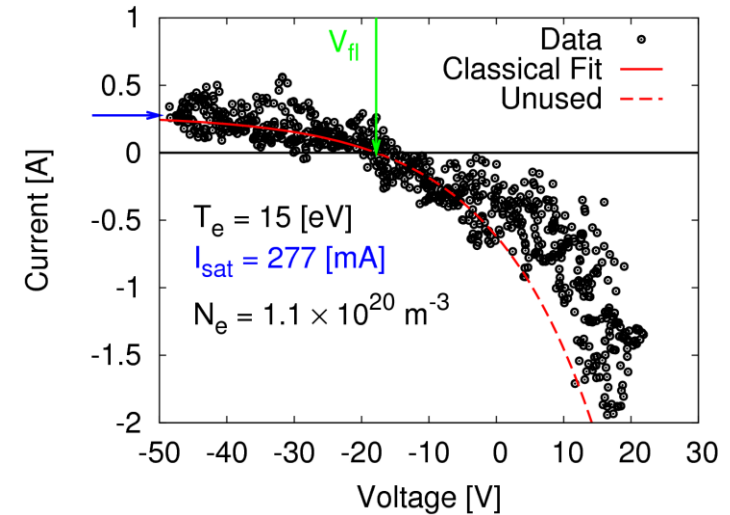
Irving Langmuir (source: wikipedia)

# Many assumptions in “text-book” theory fall down

- Typical method:
  - Subtract ion saturation; fit exponential curve to electron current
  - Derive  $N_e$  from  $T_e + I_{sat} + \text{plasma species}$
- Most derivations do not include magnetization
- Collisionality can be accounted for (see Laframboise method)
- Assumptions of Maxwellian distribution not always good to make

# Classical interpretation often yields higher temperatures relative to other diagnostics

- Classical interpretation makes use of data up to floating potential
  - Assumes single Maxwellian distribution
  - Only uses ~5% of distribution
- Independent measurements often indicate lower temperatures
  - Thomson scattering on ASDEX had some indications of non-Max. pops.
  - Thomson scattering on DIII-D consistently lower  $T_e$
  - Anomalously low sheath heat transmission coeff. on numerous machines

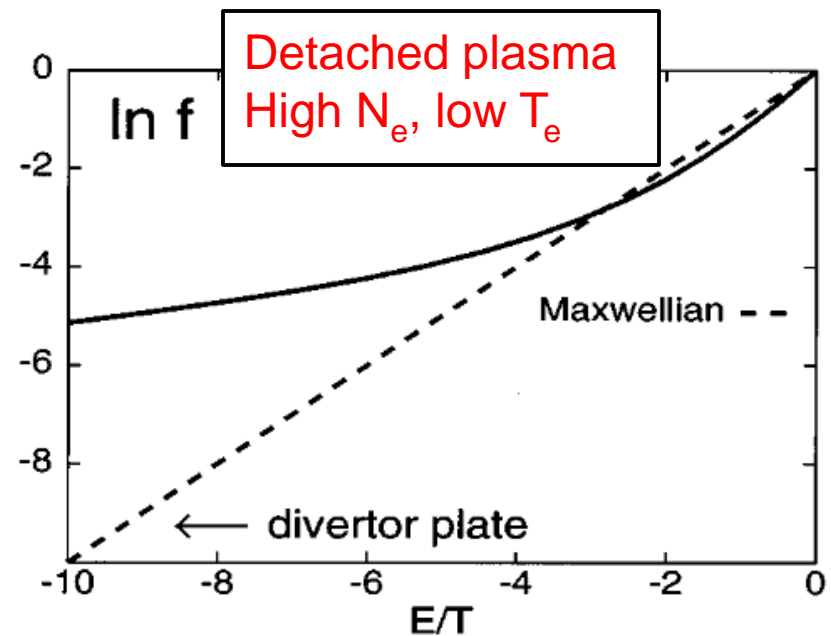


# Why expect a Maxwellian distribution?

- Maxwellian plasmas assumed due to plasma collisionality
  - Collisionality often calculated based on system length<sup>1</sup>
- Numerous modeling studies indicate non-Maxwellian distributions
  - Target plasmas result in low  $T_e$  and high  $N_e$  – yield large collisionalities in the divertor
  - Non-Maxwellian distributions still obtained
- Temperature scale length requires consideration as well

$$\lambda_{ee} \approx 1 \times 10^{16} \frac{T_e^2}{N_e} \quad \nu_{ee}^* = \frac{L}{\lambda_{ee}}$$

$L_{\parallel}$  vs.  $\lambda_{Te}$



Batischev, PoP, 1997

Other examples:

Fokker-Planck: Chodura, CPP, 1992

PIC modeling: Tskhakaya, JNM, 2011

<sup>1</sup> PC Stangeby, "The Plasma Boundary of Magnetic Fusion Devices", IoP, 2000.

# Collisionality Must Be Calculated with the Correct Scale Length

- Chodura, in 1992, pointed out the importance of local temperature scale length
- Application of Chodura criterion suggests simple limits for thermal conduction
  - Based on moments of electron distribution function
  - Most heat-carrying electrons have 3-5x thermal velocity
- **NSTX divertor plasmas indicate  $T_e > 15\text{eV}$  for fluid conduction to hold**

Temperature gradient scale length

$$L_{T_e} = \frac{T_e}{T'_e} \approx \frac{\kappa_{0e} T_e^{7/2}}{q_e}$$

Chodura criteria due to energetic electrons carrying heat

$$\frac{\lambda_{ee}}{L_{T_e}} < \frac{1}{100}$$

Scaling of minimum temperature to satisfy Chodura collisionality req.

$$T_e^{3/2} \gtrsim \frac{10^{18} q_e}{\kappa_{0e} N_e}$$

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# Kinetic probe interpretation provides more complete analysis of IV characteristics

- When electron energy scale length much longer than probe perturbation scale, velocity “diffusion” term negligible
  - $f(r,v) \rightarrow f(x,W)$
  - $W$  is total energy
  - $f_0$  is distribution far from probe
- Solution for probe characteristic determined by geometry and diffusivity
  - In magnetized plasma, cross-field diffusivity scales with Larmor radius
  - Diffusivity parameter takes form  $\psi(W) = \psi_0 W^{-1/2}$  in this case
- When  $\psi_0 \gg 1$ , first derivative becomes proportional to distribution function
- Demonstrated on CASTOR tokamak

$$x = r \quad W = \frac{1}{2}mv^2 + e\phi(x)$$

$$\lambda_\epsilon \gg r_s$$

$$\nabla_x D_x(W) \nabla_x f_0 = 0$$

$$j_e(V) = \frac{8\pi e}{3m^2\gamma} \int_{eV}^{\infty} \frac{(W - eV)f_0(W)dW}{1 + \frac{W - eV}{W}\psi(W)}$$

$$\psi(W) = \frac{1}{\gamma\lambda(W)} \int_a^{\infty} \frac{D(W)dr}{\left(\frac{r}{a}\right) D(W - e\phi)}$$

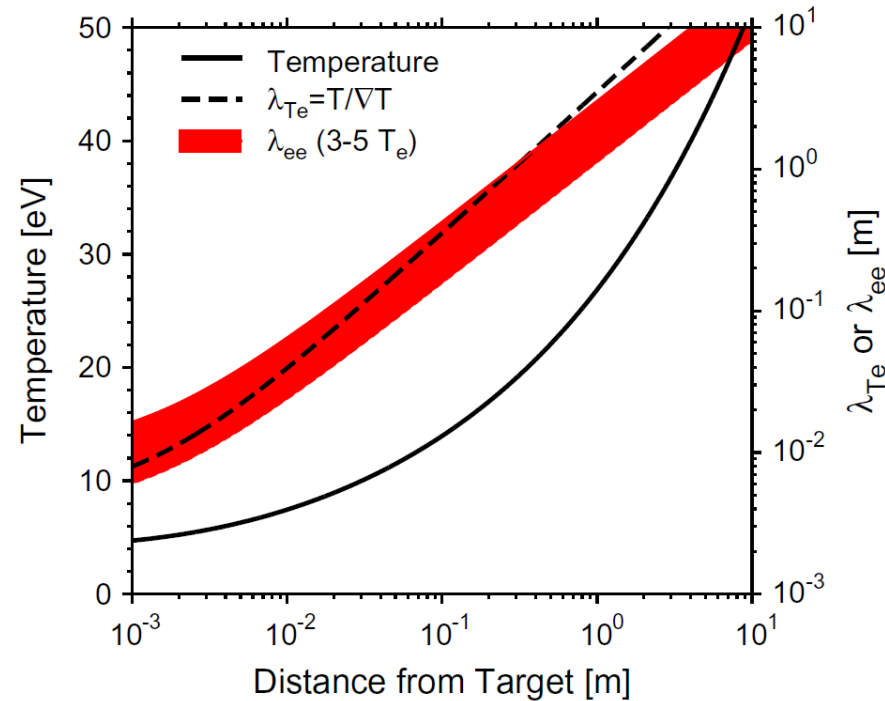
$$\boxed{\frac{dj_e(V)}{dV} \propto \frac{(eV)^{3/2}}{\psi_0} f_0(eV)}$$

Bernstein, Phys.Rev., 1954  
 Golubovskii, Sov.J.Plasma Phys, 1981  
 Arslanbekov, PSST, 1994  
 Demidov, PoP, 1999  
 Popov, PPCF, 2009  
 Godyak, Demidov, J.Phys:D, 2011

# Non-Maxwellian distributions likely: motivates kinetic Langmuir probe interpretation

- Fluid-based reconstruction (OEDGE) indicates conduction-limited regime violates Chodura conditions
- Kinetic Langmuir probe interpretation theory developed over 30 years ago
  - Golubovsky 1981 – first application to high-pressure discharges
  - Arslanbekov 1994 – application to high-pressure and magnetized discharges
  - Popov 2009 – application to tokamak edge region at midplane

• In the right conditions, **first derivative** becomes proportional to distribution function



$$\nabla_x D_x(W) \nabla_x f_0 = 0$$



$$\frac{dj_e(V)}{dV} \propto \frac{(eV)^{3/2}}{\psi_0} f_0(eV)$$

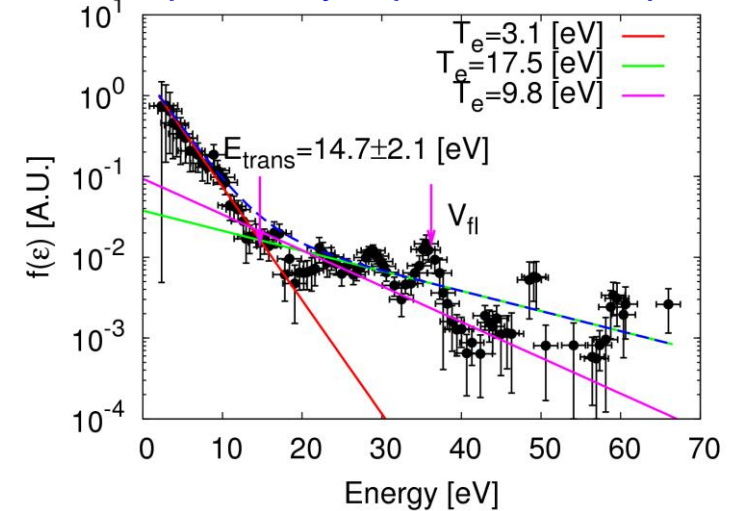
Popov, PPCF, 2009; Jaworski FED 2012, Jaworski JNM 2013



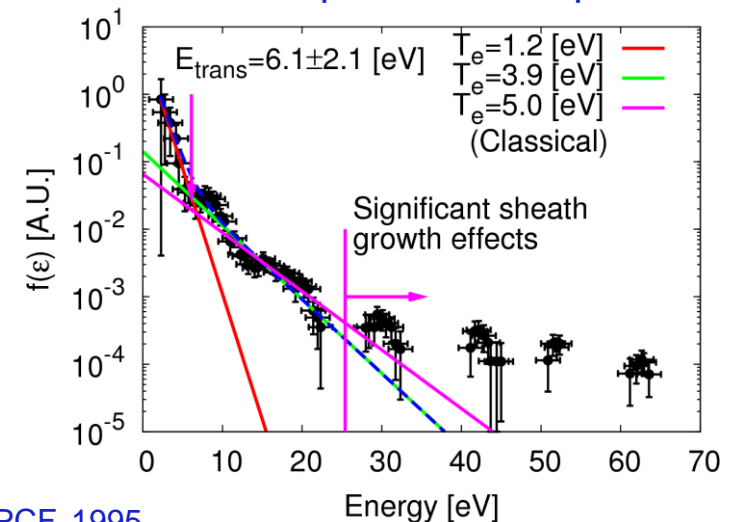
# Bi-modal distributions observed in NSTX divertor

- Typical distributions shown
  - Scrape-off layer plasma where classical  $T_e \sim 15\text{eV}$
  - Private plasma example demonstrating  $T_e \sim 1\text{eV}$
- Ion current effects due to sheath growth estimated to avoid including in fits<sup>1,2</sup>
- Bi-modal distribution often “best” model
- Total density calculated from  $I_{\text{sat}}$ 
  - Sound speed calculated using mixture of both plasma populations<sup>3</sup>

Scrape-off layer plasma example



Private plasma example

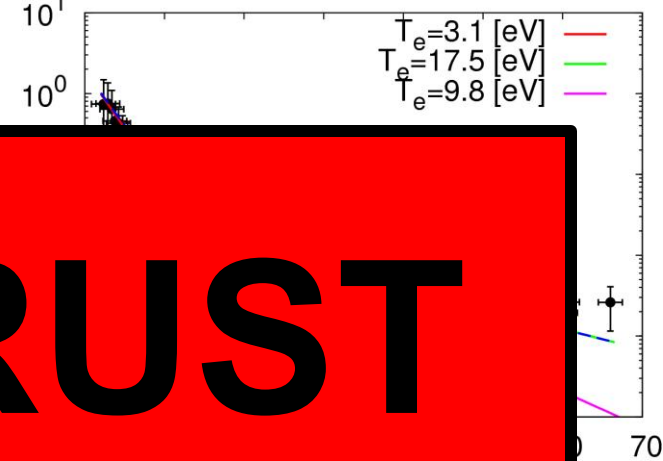


<sup>1</sup> Gunn, RSI, 1997; <sup>2</sup> Godyak, Demidov, J.Appl.Phys.D, 2011; <sup>3</sup> PC Stangeby, PPCF, 1995

# Bi-modal distributions observed in NSTX divertor

- Typical distributions shown
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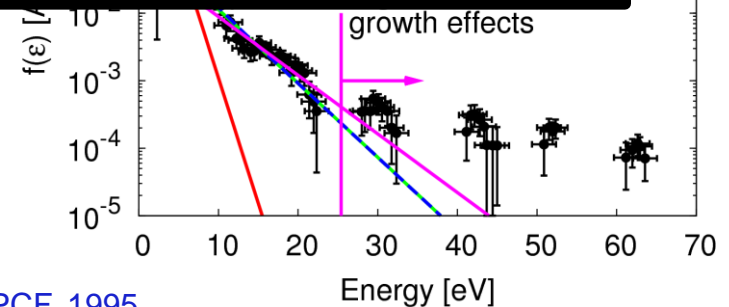
Scrape-off layer plasma example



- Ion c growth fits<sup>1,2</sup>
- Bi-modal mode
- Total

**DON'T TRUST THE PROBES!**

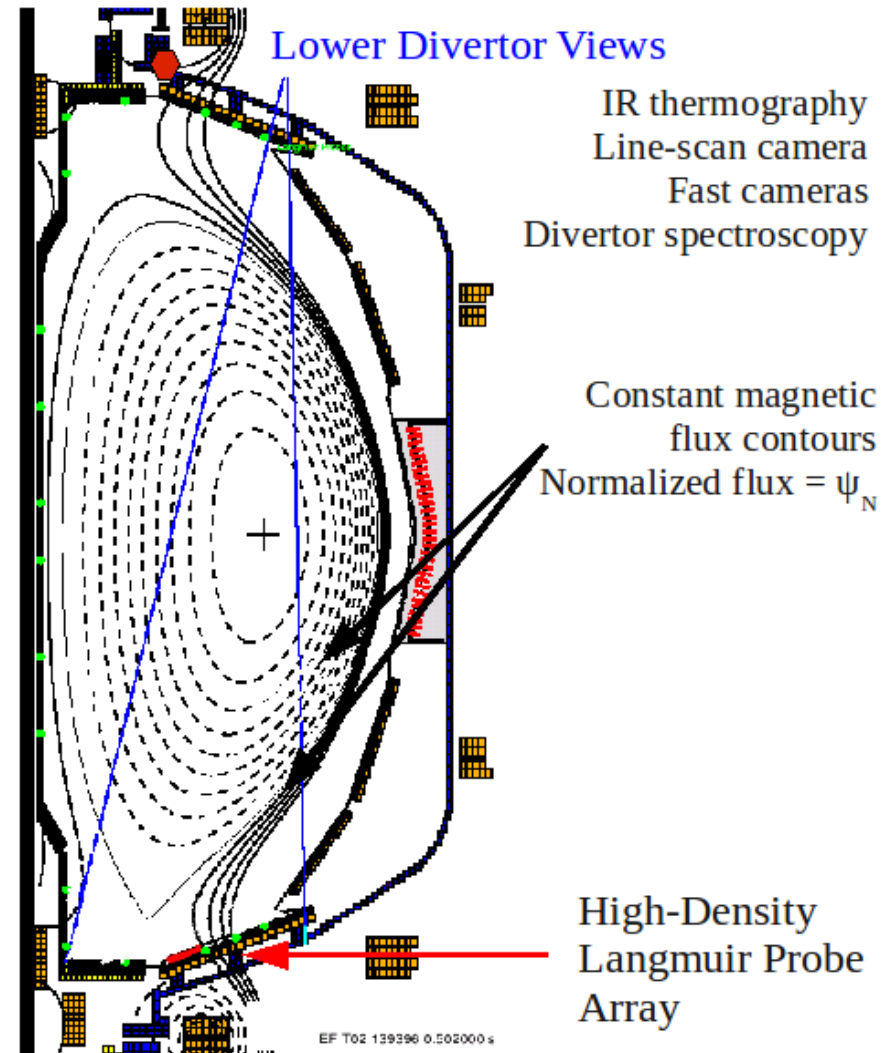
- Sound speed calculated using mixture of both plasma populations<sup>3</sup>



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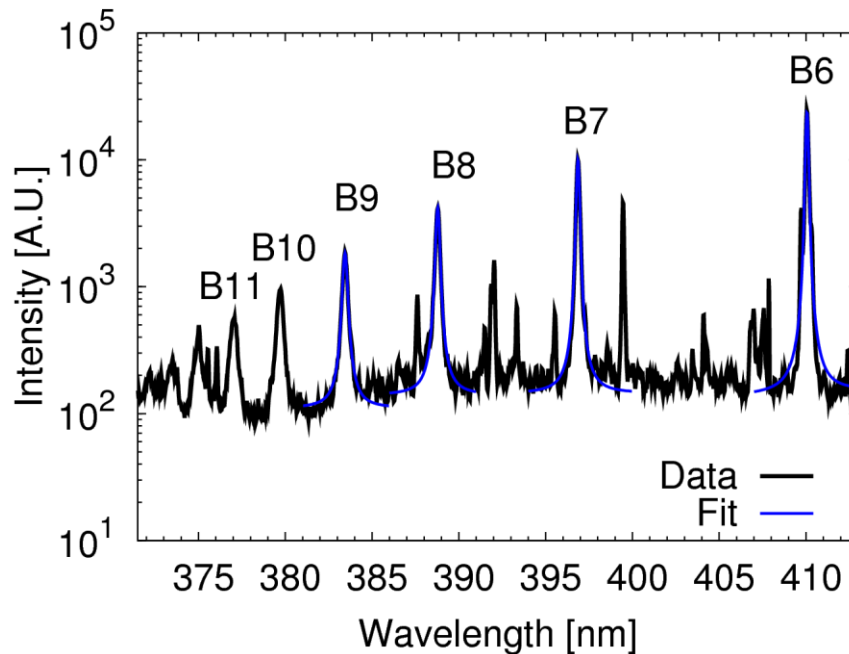
# Empirical plasma reconstruction provides framework for checking consistency between diagnostics

- Utilizes measured data points as starting point in constraining plasma models
- Solution improves as more and more data constrains background
- OEDGE code suite used here:  
Onion-Skin Method (OSM2)+EIRENE+DIVIMP
  - OSM2 fluid solver
  - EIRENE neutral hydrogen
  - DIVIMP Monte Carlo impurities
- Utilized here to provide fluid background and identify candidate diagnostics for comparison

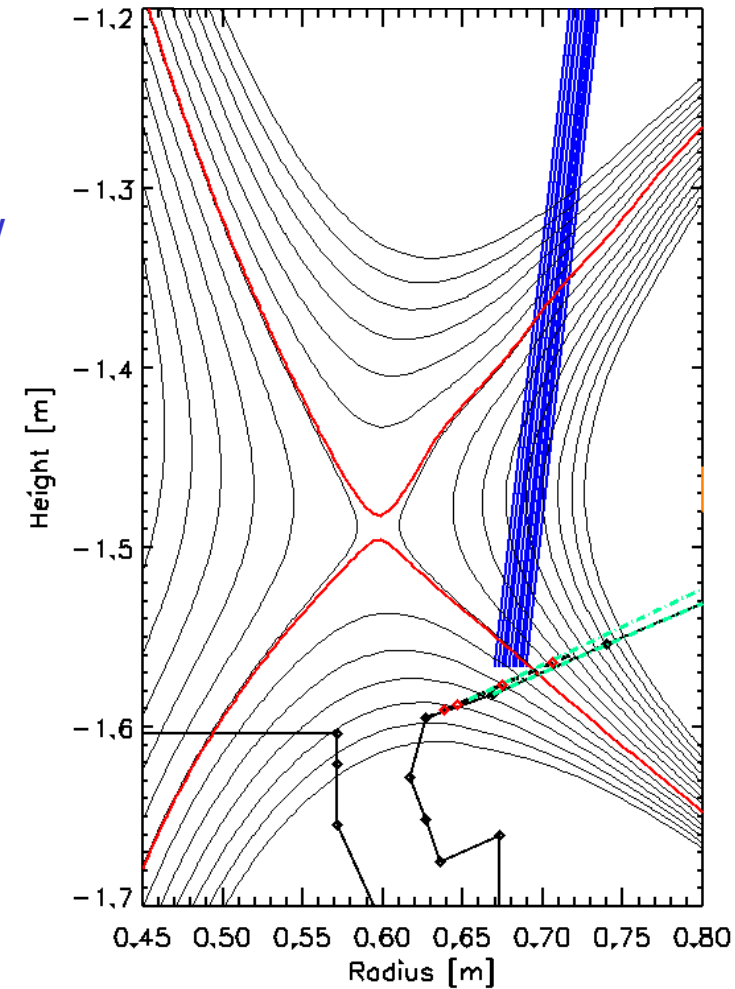


# Spectroscopy provides independent checks on density and temperature

- Divertor spectrometer viewing strike-point region during discharge
- Deuterium Balmer lines shown in this spectra
- Pressure broadening analysis indicates density of  $3.6 \times 10^{20} \text{ m}^{-3}$  (mean,  $2.1\text{-}5.5 \times 10^{20} \text{ m}^{-3}$  min/max)
  - Existence of high-n Balmer lines indicates low temperature

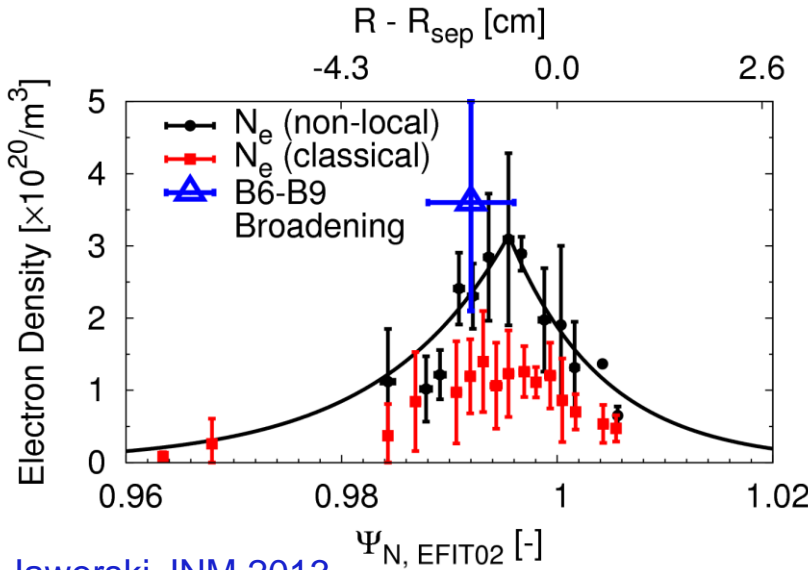


Divertor spectrometer view

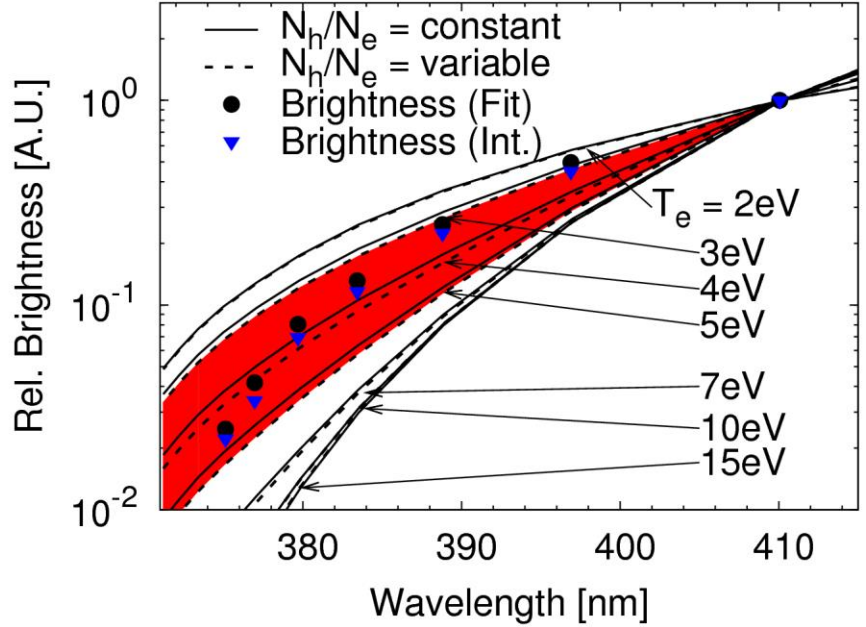
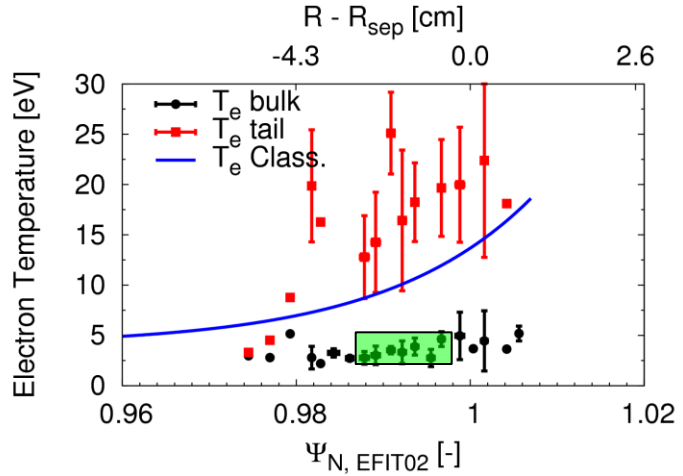


# Broadening measurement and CR modeling of hydrogen spectrum consistent with kinetic interpretation

- Pressure broadening yields density
- OEDGE plasma+neutral solution provides local parameters
- CR model calculates excited state populations for given background (Maxwellian!)
- Brightness ratios normalized to B6 consistent with 3-5eV



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# Not all models have the same knowns and unknowns!

## • Classical interpretation

- Te, Ne, and Vf come out
- Vf easily identified in I-V characteristic
- Te extracted from exponential fit
- Ne derived from Isat and Te

$$x = r \quad W = \frac{1}{2}mv^2 + e\phi(x)$$

## • Kinetic interpretation for EEDF

- Entire distribution function is up for grabs!
- Energy is referenced to **plasma potential** (not known ahead of time)
- Psi transport function not really known ahead of time (but can be guessed)

$$j_e(V) = \frac{8\pi e}{3m^2\gamma} \int_{eV}^{\infty} \frac{(W - eV)f_0(W)dW}{1 + \frac{W - eV}{W}\psi(W)}$$

$$\psi(W) = \frac{1}{\gamma\lambda(W)} \int_a^{\infty} \frac{D(W)dr}{\left(\frac{r}{a}\right) D(W - e\phi)}$$

$$\boxed{\frac{dj_e(V)}{dV} \propto \frac{(eV)^{3/2}}{\psi_0} f_0(eV)}$$

- If you don't know/trust the value, find consistency checks!

# The EEDF algorithm for NSTX-U

1. Get data from central storage
2. Determine simple IV characteristics for starting point (classical analysis)
3. Perform data smoothing (first-derivative noise)
4. Distribution function loop
  1. Construct model curve based on best guess of  $f(E)$
  2. Determine best  $V_p$  from model curve chi-square
  3. Calculate new  $f(E)$  using new  $V_p$  (also solve  $\Psi_0$ )
  4. Calculate new I-V based on  $f(E)$  and check chi-square
5. Calculate derived parameters (bi-modal  $T_e$ ,  $N_e$ , etc)
6. Write data files



# The value of chi-square (“goodness-of-fit”)

- Provides a quantitative value relating the model curve to the data set
- Not necessarily valuable as an absolute value
  - Should be  $\sim 1$  if uncertainty correctly defined
  - Less free-parameters is better
  - Uncertainty not always well defined!
  - Should be minimized for best fits
- Be wary of “black box” fitting algorithms and traps

# Consistency checks are critical!

- Make lots of plots during process development, verbose output is helpful
- Bug checking is critical – are you really applying the model you intended?
- Constantly ask yourself: is my model really better than another one?
- CHECK IT, BE QUANTITATIVE

# Other recommendations

- **COMMENT YOUR CODE**
  - You won't know when you'll get back to it or put it down
  - Comment while its fresh in your mind
- Remember what research is: you get to learn when you are wrong!

# Thank you!

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